

## Measurement of Self-Capacitance for Windings on High-Permeability Ferrite Cores

V. Yurshevich, S. Lomov and J. Jankovskis

Faculty of Electronics and Telecommunications, Riga Technical University,  
12 Azenes str., LV-1048, Riga, Latvia

**Abstract** - Two methods of self-capacitance determination for windings on high-permeability ferrite cores are described. The first one is connected with the measurements on three frequencies, the second - with the measurements on a single frequency with and without additional capacitance. The methods are free from error associated with frequency and amplitude dependencies of the complex magnetic permeability components of ferrites. The potentialities of these methods are shown by data obtained for measuring windings (single-layer windings on toroidal cores) which are broadly used in magnetic measurements. In such a complicated case when the components of complex magnetic permeability are strongly frequency dependent, the offered methods ensure the acceptable accuracy of measurements even for small self-capacitance (in order of 1 pF).

### Introduction

The self-capacitance  $C_o$  is the parasitic parameter of the windings (inductors). By increasing the frequency the influence of  $C_o$  also increases and the problem of estimation of  $C_o$  influence becomes actual. In this paper only measuring windings (MW) will be considered. MW is the single-layer winding which is wound on the toroidal core to measure its magnetic properties, for example, magnetic spectra (MS) - the frequency dependence of the components of complex magnetic permeability (CMP)  $\bar{\mu} = \mu_1 - j\mu_2$  (Fig.1). In this case the estimation of  $C_o$  influence is obligatory because the errors of MS measurements caused by inaccurate accounting of  $C_o$  influence may be comparable with errors of CMP measurement method.

Sufficiently full review of  $C_o$  measuring resonant methods is done in [1] and for other methods - in [2]. The methods described in these papers are based on two main simplifications: the inductance  $L$  and  $C_o$  (Fig. 2.a) are not changing with frequency and the inductor is linear (its parameters do not depend on the current). In many cases it is also assumed that inductor's loss are small. This makes such methods unsuitable for MW used for high-permeability ( $\mu_1 \geq 2000$ ) ferrites MS measurements. Firstly, the practice shows that in this case it is necessary to measure  $C_o$  in region II of MS (Fig. 1) where  $\mu_1$  (it means also  $L$  of MW) significantly changes with frequency. Secondly, in measuring instruments used in such methods (Q-meters, bridges) the measurement and adjustment of current through the MW, as usual, is impossible. Then, changing the frequency, changes the impedance of MW (mainly because of changes of  $\mu_1$  and  $\mu_2$ ) and consequently the current through it. The significant dependence of CMP of mentioned ferrites on the magnetic fields in specimen causes the changes of MW parameters, which may be mistakenly accredited to  $C_o$  influence. Let us consider the frequency dependence of  $C_o$ . It is known that  $C_o$  of inductors without cores may vary only in the

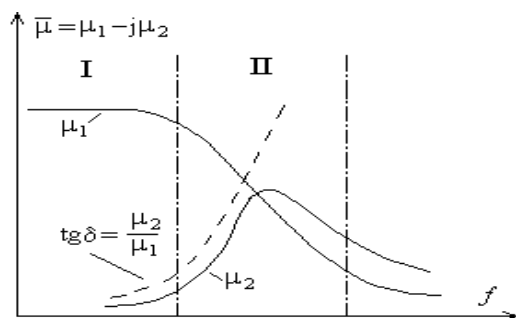


Fig. 1. Typical magnetic spectrum of high-permeability ferrite.

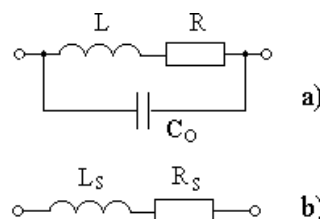


Fig. 2. Equivalent circuits of the measuring winding.

vicinity of self-resonance frequency  $f_o$ , namely at  $f/f_o > 0.6 - 0.7$ . The permittivity  $\varepsilon$  of a core must have influence on  $C_o$ . For manganese-zinc (MnZn) ferrites the value of  $\varepsilon$  is very high and it strongly changes with frequency in I and II regions of MS. However, if MW is separated from the core even with thin layer of dielectric, the influence of  $\varepsilon$  can be neglected. In this work it is accepted that  $C_o$  of MW does not depend on  $f$  right up to frequencies  $(0.6 - 0.7)f_o$ .

Thus, the authors suppose that well-known methods of  $C_o$  measurements are not valid for MW on high-permeability ferrites. In the region of low frequencies ( $f < 10 - 15$  MHz) MW has considerable advantages comparing with other transducers. Therefore, the problem of elaboration of  $C_o$  measurement methods for above-mentioned MW is important.

### First method: determination of $C_o$ from measurements on three frequencies

The essence of  $C_o$  determination methods will be explained as applied to MS measurements. To obtain the necessary formulae in such measurements usually use the equivalent circuit of MW shown in Fig. 2.a, where  $L$  is “true” inductance of MW – inductance which is free from influence of  $C_o$  as well as loss;  $R$  - resistance which is equivalent to loss in a ferrite and MW. The dielectric loss are not accounted because MW, as a rule, is separated from specimen by a good dielectric. Let us assume that the direct results of measurements at each frequency  $\omega = 2\pi f$  are the parameters  $L_S, R_S$  of series equivalent circuit of MW (Fig. 2.b). The parameters of circuits shown in Fig. 2.a and 2.b are related by well-known equations:

$$L_S = L[1 - \omega^2 LC_o(1 + tg^2\delta)] / [(1 - \omega^2 LC_o)^2 + (\omega^2 LC_o tg\delta)^2], \quad (1)$$

$$R_S = R / [(1 - \omega^2 LC_o)^2 + (\omega^2 LC_o tg\delta)^2], \text{ where } tg\delta = R/\omega L, \quad (2)$$

$$tg\delta_S = R_S/\omega L_S = tg\delta / [1 - \omega^2 LC_o(1 + tg^2\delta)]. \quad (3)$$

These equations may be used for  $C_o$  determination. After measurements on a single frequency, we obtain two equations ((1) and (2) or (1) and (3)), containing three unknown values –  $L, R$  (or  $tg\delta$ ) and  $C_o$ . Measurements on another frequency give four equations and five unknown values and so on. To obtain the determinantal simultaneous equations one needs an additional condition. As such one we can use the reality that in the case of relaxation type of MS the dependence of ferrite  $tg\delta$  from frequency is almost linear. For other MS this is valid on frequencies close to the main maximum of absorption curve  $\mu_2$ . If measurements are performed on frequencies  $f_1, f_2, f_3$  in mentioned frequency region (or in region of almost linear  $tg\delta_S(f)$  dependence), then for  $f_2 = (f_1 + f_3)/2$  we may accept that  $tg\delta_2 = (tg\delta_1 + tg\delta_3)/2$ , where  $tg\delta_i$  is the  $tg\delta$  value on frequency  $f_i$ . So we obtain the system of six equations ( $L_{S1} = \dots, tg\delta_{S1} = \dots; L_{S2}, tg\delta_{S2}; L_{S3}, tg\delta_{S3}$ ) and six unknown values ( $L_1, tg\delta_1; L_2; L_3, tg\delta_3; C_o$ ), which may be solved for  $C_o$ .

To simplify expressions we introduce such symbols:  $tg\delta_i \rightarrow t_i, tg\delta_{Si} \rightarrow t_{Si}, L_{Si} \rightarrow l_i$ . As the condition is associated with  $tg\delta$ , the following transformations were made: from equation (3) we express value  $L$ , and put it in (1). As a result we have:  $t[(1 + t_S^2)\omega^2 l C_o + 1] = t_S$ . The solution of system of three such equations:  $t_1[(1 + t_{S1}^2)\omega_1^2 l_1 C_o + 1] = t_{S1}; 0.5(t_1 + t_3)[(1 + t_{S2}^2)\omega_2^2 l_2 C_o + 1] = t_{S2}; t_3[(1 + t_{S3}^2)\omega_3^2 l_3 C_o + 1] = t_{S3}$  gives for  $C_o$  the quadratic equation:

$$C_o^2(t_{S1}A_2A_3 + t_{S3}A_1A_2 - 2t_{S2}A_1A_3) + C_o[(t_{S1}(A_2 + A_3) + t_{S3}(A_1 + A_2) - 2t_{S2}(A_1 + A_3))] + t_{S1} + t_{S3} - 2t_{S2} = 0, \text{ where } A_i = (1 + t_{Si}^2)\omega_i^2 l_i. \quad (4)$$

The solution of (4) allows to obtain  $C_o$  value. Experimental results are presented below. The weaknesses of this method are the limitation of the used frequency range for resonant MS (in this range  $tg\delta_S$  must changes with frequency almost linearly) and the influence of possible frequency dependence of  $C_o$ .

### Second method: determination of $C_o$ from measurements on single frequency

More precise results may be obtained if all measurements are carried out on the same frequency: in this case influence of possible frequency dependence of  $C_o$  is excluded. As was mentioned above, the measurements on a single frequency give two equations with three unknown values. The additional condition may be obtained if the following (second) measurements are performed on the same frequency but under different circumstances. The change of the magnetic biasing or magnetization reversal field is not rational: the corresponding relationships usually are not known precisely, but their measurements are not free from influence of  $C_o$ .

The second measurements may be done if the small capacitance  $C_{ad}$  is connected in parallel to MW. In result we obtain four equations ( $l_1, t_{S1}; l_2, t_{S2}$ ) with three unknown values ( $L, t, C_o$ ), since the parameters  $L$  and  $t$  at the same frequency are constant if  $C_{ad}$  is so small that do not change noticeably a current in the parallel circuit  $L, R, C_o$ . The presence of one “unnecessary” equation gives the opportunity to obtain the formula for  $C_o$  in different ways. One of the possible solutions:

$$C_o = [\omega^2 l_1 C_{ad} (1 + t_{S12}) + t_{S12} l_{12} - 1] / \omega^2 l_1 [1 - t_{S12} - (1 + t_{S1}^2) \omega^2 l_1 C_{ad}], \quad (5)$$

where  $t_{S12} = t_{S1}/t_{S2}$ ,  $l_{12} = l_1/l_2$ . Experimental results are presented below. The weakness of this method is difficulties with the determination of the influence of additional capacitor on  $C_o$ .

### Experimental results

It is necessary to note the contradiction arising when account of  $C_o$  influence is required. On the one hand, it is desirable to reduce  $C_o$  and, on the other hand, the error of  $C_o$  measurement sharply increases with  $C_o$  decreasing. This tendency is true for all measuring methods. For instance, in well-known method of double tuning at the resonance [1] at frequency ratio  $n=4$  the relative error  $\Delta C_o/C_o$  increases from  $\pm 28\%$  for  $C_o = 5 \text{ pF}$  to  $\pm 270\%$  for  $C_o = 0.5 \text{ pF}$ . As practice shows, the ordinary MW for high-permeability ferrites MS measurements have  $C_o$  values in range  $0.5 - 2 \text{ pF}$ .

The possibilities of the first method may be shown by using the results of  $C_o$  measurements for MW wound on the toroidal cores of MnZn ferrite with initial magnetic permeability  $\mu_a = 3000$ ; the dimensions of toroidal core are  $20 \times 12 \times 6 \text{ mm}$ ; number of turns  $w = 24$ . The preliminary measurements of MS showed that dependence  $tg\delta_s(f)$  is almost linear in range  $0.7 - 1.7 \text{ MHz}$ . In this range  $C_o$  was determined using (4) for three similar specimens and several frequency sets, for example, from  $(0.7 - 0.85 - 1.0) \text{ MHz}$  till  $(1.3 - 1.5 - 1.7) \text{ MHz}$ . As it was expected, the scattering of obtained  $C_o$  values is high enough and increases with decreasing the frequency (decreasing of the ratio  $f/f_o$ ). For frequency interval  $1.2 - 1.7 \text{ MHz}$  the calculated values of  $C_o$  were located in the range from  $0.3$  to  $1.6 \text{ pF}$  but the great part of results were concentrated about mean value  $C_{om} = 0.95 \text{ pF}$  at scattering  $\pm 0.4 \text{ pF}$ . It should be pointed out that the mentioned scattering is not the weakness of the method. This is the objective phenomenon: the precise measurements of small  $C_o$  at the limited value of  $f/f_o$  is a difficult problem (the influence of  $C_o$  is slight and the relationship between the effective parameters is complicated) Thus, the elements of statistical approach are natural for such measurements.

For above-mentioned MW  $C_o$  were measured on frequencies  $0.6$  and  $0.75 \text{ MHz}$  also by second method at  $C_{ad} = 2.4$  and  $4.1 \text{ pF}$ . It is obtained that  $C_{om} = 1.1 \text{ pF}$  at real scattering  $\pm 0.4 \text{ pF}$ . The rise of  $C_{om}$  may be explained by increase of current in parallel circuit  $L, R, C_o$  (Fig. 2.a) after connection of  $C_{ad}$  (that is at approaching to self-resonance). In this case according to amplitude dependence of  $\mu_1$  the inductance of MW increases and this change, in fact, is ascribed to influence of  $C_o$ .

For MW on the same cores but with  $w = 72$  the self-capacitance has been measured at the same frequencies and the same  $C_{ad}$ . It was obtained that  $C_o = 1.4 \pm 0.3 \text{ pF}$ . At the three-fold increasing of

with the self-capacitance increases only for 40%. This means that contribution to  $C_o$  from MW leads and parasitic capacitance of measuring network is noticeable.

In this study the results of MS measurements were used for  $C_o$  determination by the first method. If  $C_o$  determination is not associated with MS measurements, one must use the measuring method which allows to monitor and adjust the current through MW (at MS measurements this is an obligatory condition), otherwise the amplitude dependence of complex magnetic permeability may introduce the uncertainty in the results. Besides that such a method allows to correct an increase of the current in parallel circuit L, R,  $C_o$  when  $C_{ad}$  is connected.

The great scattering of results and its increase with decreasing of  $f$  may be explained easily on the

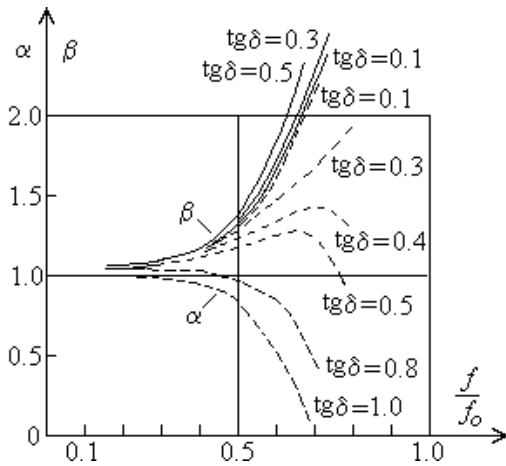


Fig. 3. Frequency dependence of parameters  $\alpha, \beta$ ;  $\alpha = L_S/L$ ;  $\beta = tg\delta_S/tg\delta$ .

basis of Fig.3: when loss and  $f/f_o$  are small, the parameters  $L_S$  and  $tg\delta_S$  are changing insignificantly and almost in the same extent. Then the equations (3) and (4) contain the differences between the closely spaced quantities and even small errors of  $L_S$  and  $tg\delta_S$  determination involve the significant scattering of  $C_o$  values. From this the importance of the increasing of measurement accuracy follows (the individual calibration of measuring instruments, the accounting of main accompanying factors: the parameters of instruments' input, the parasitics of measuring network a. o.). For example, even after simplest performance of these operations the scattering of  $C_o$  decreased to  $\pm 0.25$  pF for MW with  $w = 24$ .

### Conclusion

The offered methods allow to measure with acceptable accuracy even small  $C_o$  on frequencies where the components of complex magnetic permeability of ferrites are changing strongly. The well-known methods are unsuitable for this case as they are not taking into account these changes and, in fact, the corresponding changes of MW  $L_S$  and  $tg\delta_S$  are accredited to influence of self-capacitance. As it was mentioned, the results of MS measurements were used for  $C_o$  determination by the first method. This is undoubted advantage of the method if  $C_o$  values are necessary for correction of MS measurements. If the measurements of  $C_o$  are not associated with MS, the allowable range of frequencies may be easily found from the frequency dependence of parameter  $tg\delta/\mu_a$ . One can find these curves in handbooks. The weakness of the second method is connected with the using of additional capacitor. However, the comparison of the results obtained by both methods showed that this influence is insignificant if one uses the miniature capacitors and  $f/f_o < 0.6$ .

These methods were developed mainly for MS measurements. But they conceptually are valid also for other windings on high-permeability ferrite cores – for windings of transformers and chokes.

### References

1. J. P. Newsome, *Electronic Engineering*, **28**, august, 350 (1956).
2. V. Yurshevich, I. Presnova, in *Problems of electrodynamics and of electronic circuits theory* (Riga, Latvia), 35 (1980) (in Russian).