

Measurement of the bandwidth of Helicon Maser

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***Abstract.** Radio frequency (RF) magnetoplasmic waves known as helicons will propagate in solid - state plasmas of metals and semiconductors when a strong magnetic field is applied. In Helicon Maser (HM) the helicons are excited by RFs much higher than the helicon generation frequency. The higher frequency e/m field may be called a pumping field. In analogy with the usual maser (or laser) the magnetized semiconductor sample plays the role of active material and the connecting cable – the role of high quality external resonator. The bandwidth of HM is much narrower than in case of isolated sample (the same situation as in case of laser). The new method of HM bandwidth measurement is considered. InSb and Ge were used for active materials. The constant magnetic field varied within the interval 0 – 23 Tesla.*

Keywords: Helicon Maser; Magnetoplasmic waves; Microwave resonator.

1. Introduction

Helicons are magnetoplasmic waves propagating in semiconductors and metals when a strong magnetic field is applied. In our paper [1] the construction of Helicon Maser (HM) was described. The generation of helicons in the HM was obtained by the help of maser effect when the magnetoplasmic waves were excited by radio frequencies (the range 1000 MHz) much higher than the helicon generation frequency (main peak at 20 MHz). The higher frequency field may be called a pumping field. In analogy with the usual maser (or laser) the semiconductor sample plays the role of an active material and the connecting cable – the role of high quality external resonator. The bandwidth of the HM is much narrower than the bandwidth of isolated sample (the same situation as in case of laser).

The HM may be used for the contactless measurement of carrier density and mobility in semiconductors and semiconductors, RF antenna modelling, investigation of self – focusing and diffraction of e/m waves, etc [2].

The measurement of the HM bandwidth in working mode is rather complicated as the duration of helicon impulse is about 10 msec. In this paper the new experimental method of the HM bandwidth measurement is described.

2. Subject and Methods

The block diagram of the Helicon Maser is shown in Fig.1. Here 1 is semiconductor sample in the form of plate; 2 is excitation coil; 3 is pickup coil; 4 and 5 are the coaxial cable lines.

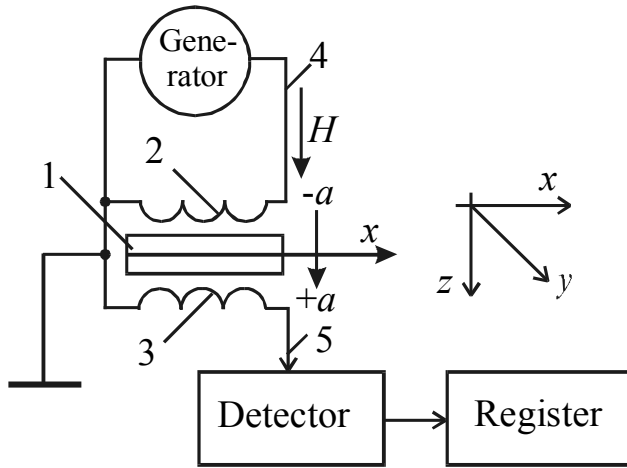


Fig. 1. The experimental device of Helicon Maser

The direction of the coordinate axes x , y , z are also shown. The constant magnetic field H is directed along the axis z and is perpendicular to the sample plane x - y . The sample thickness is $2a$ and the dimensions of the plate in x - y direction is much larger. In this case excited helicon is a circularly polarised plane wave of the form $\exp [i(kz-\omega t)]$.

The propagation vector k is taken to be in the z direction and may be calculated from the dispersion relationship (it is assumed that $\omega_H \tau \gg 1$)

$$k = \sqrt{\frac{\omega_p^2 \cdot \omega}{\omega_H \cdot c^2} \left(1 + \frac{i}{2\omega_H \tau} \right)}, \quad (1)$$

where ω_p is the plasma frequency, ω_H is the cyclotron frequency, τ is the collision time, c is the velocity of light [3].

For the n -type semiconductor the plasma and cyclotron frequencies are given by

$$\omega_p^2 = \frac{4\pi N e^2}{m}; \quad \omega_H = \frac{eH}{mc}, \quad (2)$$

where N , e , m are the density, charge and effective mass of the electrons (Gauss units). The resonance frequency of the active material (semiconductor) may be found from the condition

$$(\text{Re } k)2a = \frac{n\pi}{2}, \quad n = 1, 3, 5, \dots \quad (3)$$

The Helicon Maser instabilities (generation and/or amplification) occur when the resonant frequency of the active material coincides with the resonant frequency of the cable line.

The quality of the semiconductor sample (active material) $Q_1 = \omega_H \tau \leq 10$. The long cable line quality Q_2 in the resonant mode is much higher $Q_2 = 4000$. The equivalent quality of HR Maser (active sample + long line resonator) in our experiments reaches the value $Q = 2000$.

The most appropriate material for the experimental realization of HM is n -InSb where $\omega_H \tau = 1$ condition may be achieved in the fields of 0.2 T (at the room temperature).

In the helicon maser (HM) is connected to the pumping RF generator with the frequency $f_2 = 100 - 2000$ MHz much higher than the helicon resonance frequency $f_1 = 20$ MHz the pick-up signal during the generation time is proportional to time function

$$F = \frac{f_2}{f_1} \sin 2\pi f_1 t - \sin 2\pi f_2 t \quad (4)$$

with the initial condition

$$F(0) = F'(0) = 0. \quad (5)$$

3. Experimental results

During the helicon generation time the first item in formula (4) represents the Helicon Maser response and is much larger than the second item (the pumping signal). Experimental results were obtained for a sample n-InSb (carrier density $N = 1,5 \cdot 10^{16} \text{ cm}^{-3}$; mobility $\mu = 5 \cdot 10^4 \text{ cm}^2 \text{ V}^{-1} \cdot \text{S}^{-1}$; sample thickness $2a = 0,5 \text{ cm}$; $T = 300 \text{ K}$). The main helicon generation peak was observed at $B = 1,55 \text{ T}$ with the helicon frequency $f_1 = 20 \text{ MHz}$ in the wide range of pumping frequencies $f_2 = 100 \div 2000 \text{ MHz}$. The typical Helicon Maser response is shown in Fig. 2 where the lower generation frequency f_1 peaks are superimposed on the higher frequency f_2 of the pumping field (eq.(4)).

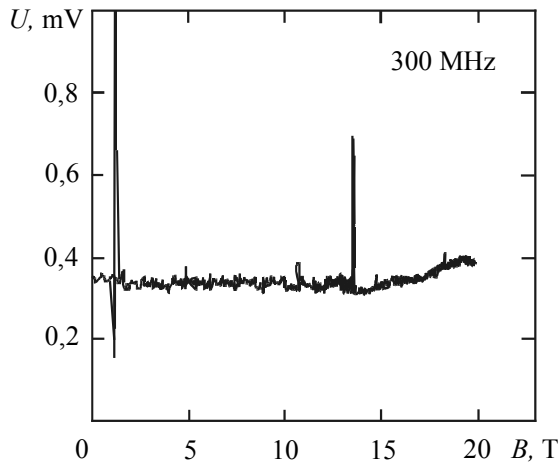


Fig. 2. The pick-up signal in millivolts versus constant magnetic field induction B in Tesla (pumping 300 MHz)

4. The HM bandwidth measurement results

The direct measurement of the HM bandwidth in working mode is rather complicated as the duration of helicon impulse is only about 10 msec. In this paper the new measurement method is described and all operations may be provided without pumping field.

The excitation coil in this case is connected to the RF source of the same frequency as helicon generation. The process of the measurement is illustrated by Fig. 3, where U is the signal in the pickup coil and B is induction of constant magnetic field.

In experimental device Fig. 1 the generation frequency is 20 MHz. If the HM is excited by the RFs slightly less than 20 MHz we have the dependence $U = f(B)$ described by the curves 1, 2 and 3. The resonant curve 4 for $f_1 = 20 \text{ MHz}$ has the smallest bandwidth which coincides with the bandwidth of Helicon Maser in the working mode when pumping field (100 – 2000 MHz) is applied. For slightly higher excitation RF we have once again the wider curve 5. Experimentally the helicon generation frequency and HM bandwidth may be determined by starting at the lower RF (curve 1) and continuously increasing the initial frequency.

In all cases the carrier density in the active material may be found by fixing the generation frequency f_1 and using the equations (1), (2) and (3) where $\omega = 2\pi f_1$.

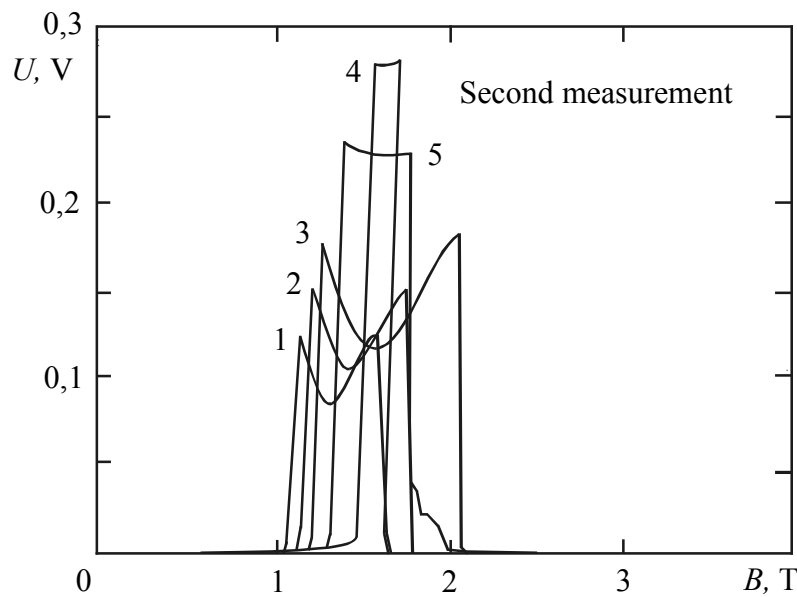


Fig. 3. The resonance curves of the HM for various frequencies of excitation. 1 – 19,3 MHz; 2 – 19,5 MHz; 3 – 19,7 MHz; 4 – 20 MHz (generation frequency); 5 – 20,5 MHz; U – the signal of pickup coil in Volts; B – magnetic field induction in Tesla

5. Conclusions

In the Helicon Maser magnetoplasma waves are excited by the radio frequencies much higher than the helicon generation frequency. The excitation of helicons in this case may be described by the effect similar to the Combination Scattering (Raman effect) when a part of the high RF wave energy that passes through the active material is absorbed and re-emitted by the magnetized solid – state plasma. The high RF e/m field may be called the pumping field. In analogy with the usual maser (or laser) the semiconductor sample plays the role of the active material and the connecting cable – the role of high quality external resonator. The bandwidth of the HM is much narrower than the bandwidth of isolated sample (exactly the same situation as in case of laser). The direct measurement of the HM bandwidth in working mode is complicated as the duration of helicon impulse is only about 10 msec. The new measurement method is proposed. Experimental results for an active material n-InSb were obtained in the range of constant magnetic fields $B = 0 - 20$ T and pumping frequencies 100 – 2000 MHz.

References

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